Mini-Review on Mini-Black Holes from the Mini-Big Bang



Marcus Bleicher , ITP, Goethe Universität Frankfurt, Germany





relativistic heavy ion collider

anterline of 010

Statement on Committee Review of Speculative "Disaster Scenarios" at Brookhaven Lab's RHIC

October 6, 1999

Brookhaven National Laboratory has posted on its Web site a report by expert physicists who recently reviewed speculative disaster scenarios at the Relativistic Heavy Ion Collider.

The report summarizes technical discussions that conclude there is no danger of a "disaster" at RHIC

In July 1999, Brookhaven Lab Director John Marburger convened a committee of distinguished physicists to write a comprehensive report on the arguments that address the safety of each of the speculative disaster scenarios at RHIC. The scenarios are:

- Creation of a black hole that would "eat" ordinary matter.

- Initiation of a transition to a new, more stable universe.

triggered more than 1000 par and a couple of lawsuits)

Formation of a "strangelet" that would convert ordinary matter to a new form.

"We conclude that there are no credible mechanisms for catastrophic scenarios at RHIC," said committee chair Robert Jaffe, Professor of Physics and Director, Center for Theoretical Physics at Massachusetts Institute of Technology. "Accordingly, we see no reason to delay RHIC operation."

Thanks to



- Ben Koch (ITP, Germany)
- Ulrich Harbach (\rightarrow Patent lawyer)
- Christoph Rahmede (Sissa, Italy)
- Piero Nicolini (FIAS, Germany)
- Martin Sprenger (DESY, Germany)
- Sabine Hossenfelder (NORDITA, Stockholm)
- Stefan Hofmann (LMU Munich, Germany)

Based on: Phys.Rev.D66:101502,2002, Phys.Lett.B548:73-76,2002, J.Phys.G28:1657-1665,2002, Phys.Lett.B566:233-239,2003, Int.J.Mod.Phys.D13:1453-1460,2004, JHEP 0510:053,2005





Happy Birthday Takeshi!



The Standard Model



Carrier	Force	Group	
γ photon	E&M	U(1)	
<i>g</i> gluon	Strong	SU(3)	
Z W	Weak	SU(2)	

Fermilab 95-759







- Essential result: Me
- More dimensionsNoncommutative space

Motivation for extra dimensions



- Kaluza-Klein (last century 1920ies)
- String theory/M-theory/SUGRA
 →10/11 dimensions
- Hierarchy problem can be solved (why is the Planck scale so much different from the electro-weak scale?!)
- Usually compactified extra-dimensions on radii R~1/M_{Planck} (hard to observe)

Extra dimensions



- Arkani-Hamed, Dvali, Dimopoulos (1998)
 - \rightarrow 1-7 large extra dimensions

$$\rightarrow$$
 only for gravity

- \rightarrow R~ fm ... mm (y_i are compactified)
- → new fundamental scale M_f~TeV

$$(x_0, x_1, x_2, x_3, y_1, \dots, y_d) =: (x, y)$$

$$ds^{2} = -(dx^{0})^{2} + \sum_{\nu} (dx^{\nu})^{2} + \sum_{i} (dy^{i})^{2}$$

N. Arkani-Hamed, S. Dimopoulos and G. R. Dvali, PLB **429**, 263 (1998);

See also RS model \rightarrow 1 warped extra dim. R~10 I_Planck

Einstein-Hilbert action

- Expansion of the metric (KK-towers)
- Integration of the d spatial dimensions

$$S_{d+4} = M_f^{d+2} \int R \sqrt{g} \mathrm{d}^4 x \mathrm{d}^d y$$

= $L^d M_f^{d+2} \int {}^{(4)} R \sqrt{{}^{(4)}g} \mathrm{d}^4 x + \sum_{n>0} (\dots)$
 $\Rightarrow m_p^2 = M_f^{d+2} L^d$



Matching Newton 's law

Newton with **LXDs**:



Observables of LXDs





Tests of Newtonian Gravity









Black Holes



- Solutions to Einstein's equations
- Schwarzschild radius r_s ~ GM_{BH} – requires large mass/energy in small volume
- Light and other particles do not escape; classically, BHs are stable



The Schwarzschild radius



 $ds^{2} = -\alpha_{(d+3)}(r,t)dt^{2} + \beta_{(d+3)}(r,t)dr^{2} + r^{2}d\Omega_{(d+3)}$

Myers, Perry 1984

D-dimensional Einstein Eq. gives the **Schwarzschild radius** for a D-dimensional Black Hole

$$R_S^{d+1} = \frac{2}{d+1} \left(\frac{1}{M_f}\right)^{d+1} \frac{M}{M_f} \qquad \qquad \text{Surface} \qquad \kappa_D = \frac{1+d}{2} \frac{1}{R_S}$$

I.e. $R_s \sim M/m_{grav}^2$ Reduction of $m_{grav} \rightarrow increase in R_s$



• BH cross section is

 $\rightarrow \sigma(p_1 p_2 \rightarrow BH) \sim 400 \text{ pb}$



Black Holes at Colliders

• BH created when two particles of high enough energy pass within $\sim r_s$.

Eardley, Giddings, PRD (2002) Yoshino, Nambu, PRD (2003) Dimopoulos, Landsberg, PRL (2001)

- Large Hadron Collider: $E_{COM} = 14 \text{ TeV}$ $pp \rightarrow BH + X$
- LHC prediction: 100
 black holes per second

$$\frac{\mathrm{d}\sigma}{\mathrm{d}M} = \sum_{A_1,B_2} \int_0^1 \mathrm{d}x_1 \frac{2\sqrt{\hat{s}}}{x_1s} f_A(x_1,\hat{s}) f_B(y_2,\hat{s}) \sigma(\hat{s},d)$$



LHC: S. Hossenfelder, M.B., et al., PRD 66 (2002) Tevatron: M. Bleicher et al., PLB548 (2002)

Signatures



- Event shapes (thermal emission)
- Modifies hadron spectra
- Exotic particle production
- Cut-off in pT spectra
- Remnants?

Constraints from LHC



- Exotic particle search
- Constraints on the exchange of gravitons
- Kaluza-Klein excitations
- Direct searches for black holes



Graviton exchange



Experiment	Process	+	_
LEP [7]	$e^+e^- ightarrow \gamma\gamma$	$0.93{ m TeV}$	$1.01{ m TeV}$
LEP $[8-11]$	$e^+e^- \rightarrow e^+e^-$	$1.18\mathrm{TeV}$	$1.17{ m TeV}$
H1 [12]	e^+p and e^-p	$0.74{ m TeV}$	$0.71{\rm TeV}$
ZEUS $[13]$	e^+p and e^-p	$0.72{ m TeV}$	$0.73{ m TeV}$
CDF [14]	$p\bar{p} \rightarrow e^+e^-, \gamma\gamma$	$0.99{\rm TeV}$	$0.96{\rm TeV}$
DO[14]	$p\bar{p} \rightarrow e^+e^-, \gamma\gamma$	$1.28{ m TeV}$	$1.14\mathrm{TeV}$
DO[15]	$p\bar{p} ightarrow jj$	$1.48 { m TeV}$	$1.48{ m TeV}$
CMS at 7 TeV with $40/\text{pb}$ [16]	$pp \rightarrow \mu^- \mu^+$	$1.6{ m TeV}$	$1.6{ m TeV}$
CMS at 7 TeV with $36/\text{pb}$ [30]	$pp \rightarrow \gamma \gamma$	$1.74\mathrm{TeV}$	$1.71{ m TeV}$
ATLAS at $7 \mathrm{TeV}$ with $3.1/\mathrm{pb}$	pp ightarrow jj	$2.2{ m TeV}$	$2.1{ m TeV}$
ATLAS at $7 \mathrm{TeV}$ with $36/\mathrm{pb}$	pp ightarrow jj	$4.2{ m TeV}$	$3.2{ m TeV}$
CMS at $7\mathrm{TeV}$ with $36/\mathrm{pb}$	pp ightarrow jj	$4.2{ m TeV}$	$3.4{ m TeV}$



Kaluza-Klein excitations



- 3+d space like dimensions
- d dimensions on d-torus with radii R
- only gravity propagates in all dimensions (bulk)
- all other in 4-dim. space time (brane)

$$^{(5)}\Box\Phi = 0$$

$$\Leftrightarrow \left({}^{(4)}\Box + \partial_5^2 \right) \Phi = 0$$

$$\Leftrightarrow \left({}^{(4)}\Box - \frac{n^2}{R^2} \right) \Phi = 0$$





Randall-Sundrum KK (II)



I-Surandall-Surgravitohs





(fb)

from Evan Wulf (ATLAS)





2011 $G \rightarrow ee restriction Randall-Sundrum KK (IV)$





Black Hole Evaporation

- Quantum mechanically, black holes are not stable – they emit Hawking radiation
- Temperature: $T_{\rm H} \sim 1/{\rm R_s}$ Lifetime: $\tau \sim (M_{\rm BH})^3$
- For M_{BH} ~ M_{sun}, T_H ~ 0.01 K. Astrophysical BHs emit only photons, live ~ forever







Exotic particle production





Direct black hole search



9 Jet event, S_T=2.6 TeV

CMS Experiment at LHC, CERN Data recorded, Mon May 23 21.46-26 2011 EDT Run/Event 156567 / 347495624 Lumi section: 200 Orb/Crossing: 73256853 / 3161



Event Characteristics

Decay Signature Average of ~ 6 particles for each decay, emitted spherically ~120 Particle degrees of freedom

> Summing over spin and color gives: 75 % quarks and gluons 10 % charged leptons 5 % neutrinos 5 % photons or W/Z bosons Also get new particles around 100-1000 GeV, Small fraction of invisible neutrinos and gravitons → BH's easy to reconstruct

10% high PT leptons \rightarrow trigger



De Roeck (2002) S. Vahsen (2008) Marcus Bleicher, RANP 2013

Hadron spectrum



Get hadron spectrum from parton fragmentation

$$E\frac{d\sigma^{h}}{d^{3}p} = \frac{1}{s} \sum_{a,b,c} \int_{M_{BH,min}^{2}}^{s} dM_{BH}^{2} \int_{x_{1,min}}^{1} \frac{dx_{1}}{x_{1}} \int_{z_{min}}^{1} \frac{dz}{z^{2}}$$

× $f_{a}(x_{1},Q^{2})\sigma_{BH}f_{b}(x_{2},Q^{2})E_{c}\frac{dN_{c}}{d^{3}p_{c}}D_{c}^{h}(z,Q_{f}^{2}),$

- Charged hadrons from BHs exceed pQCD at high pT
- Bump near Hawking temperature



Direct black hole search (I)





- Search for high E_T multi-particle final states
- Similar results for N>2 final state particles
- No deviations from Standard Model observed!



CMS, arxiv:1012.3375

M_{BH}^{min} (TeV)

Direct black hole search (II)





- 95% percent exclusion limits
- All cases (rotating, non rotating, remnant formation) are strongly constrained



CMS, arxiv:1012.3375

Conclusion



- Mini Black Holes, Gravitons, extra dimensions are still an interesting idea
- However:
 - No signals of BHs observed at LHC
 - Also no signal of Gravitons at LHC
 - (Also no signal of non-commutativity)







Remnants



- Include remnant in Charybdis by a modification of the emission spectrum
- Try direct measurement of heavy charged remnant
- \rightarrow Ashes of the black hole

See also Bonano, Reuter (Renormalized coupling constant) and Rizzo et al (Modified gravity)

→ Same effect, different origin



Koch, Hossenfelder, Bleicher (2007)



Black Holes from Cosmic Ray

- Cosmic rays Nature's free collider
- Observed events with 10^{20} eV produces $E_{\text{COM}} \sim 500 \text{ TeV}$
- But meager fluxes. Can one use this energy?



Auger Observatory





Black holes from cosmic neutrinos

- Currently no such events seen → stringent bound on extra dimensions.
- Auger can detect
 ~100 black holes in
 3 years.





Feng, Shapere, PRL (2002



AMANDA/IceCube

- Neutrino telescopes: promising BH detectors
- Similar rate: ~10 BH/year
- Complementary information
 - BH branching ratios (jets vs. muons)
 - Angular distributions

Kowalski, Ringwald, Tu, PLB (2002) Alvarez-Muniz, Feng, Han, Hooper, Halzen, PRD (2002) Harbach, Bleicher, subm. Astroparticle Phys.(2005)









Luminosity difference: 10²⁷ vs 10³⁴ /cm/s

Conclusions



- Gravity is either weak or is strong but diluted by extra dimensions
- Black hole production is a leading test



 If gravity is strong at the TeV scale, we will find microscopic black holes at LHC!

Black Hole evaporation?

- "Balding phase":BH gets rid of its hair mainly via gravitational radiation ->not visible in detector
- Hawking phase: decay mainly into standard-model particles.
 - R. Emparan, G. T. Horowitz and R. C. Myers Phys. Rev. Lett. 85, 499 (2000)
 - S. B. Giddings and S. Thomas, Phys. Rev. **D 65** 056010 (2002).
 - C. M. Harris, M. J. Palmer, M. A. Parker, P. Richardson, A. Sabetfakhri and B. R. Webber, [arXiv:hep-ph/0411022]...
- Final state: Two possible scenarios
 - Hawking radiation continues until M_{BH}....M_f and then performes something like a final decay
 - Rapid decay slows down to form quasi-stable remnant
 - Y. B. Zel' dovich, in: "Proc. 2nd Seminar in Quantum Gravity", edited by M. A. Markov and P. C. West, Plenum, New York (1984).
 - R. J. Adler, P. Chen and D. I. Santiago, Gen. Rel. Grav. 33, 2101 (2001)
 - J. D. Barrow, E. J. Copeland and A. R. Liddle, Phys. Rev. **D 46**, 645 (1992).
 - S. Coleman, J. Preskill and F. Wilczek, Mod. Phys. Lett. A6 1631 (1991).
 - S. Hossenfelder, M. Bleicher, S. Hofmann, H. Stocker and A. Kotwal, Phys. Lett. B 566, 233 (2002)
 - M. Bonanno, M. Reuter, Phys. Rev. D 73 (2006)





 $T_H = \frac{d+1}{4\pi R}$

Still vital discussion of this (classical) cross sections



- S. Dimopoulos and G. Landsberg Phys. Rev. Lett. 87, 161602 (2001).
- M. B. Voloshin, Phys. Lett. B **518**, 137 (2001); Phys. Lett. B **524**, 376 (2002).
- S. B. Giddings, ed. N. Graf, eConf C010630, P328 (2001).
- S. N. Solodukhin, Phys. Lett. B 533, 153 (2002).
- H. Yoshino and V. S. Rychkov, Phys. Rev. D 71 (2005) 104028 ...

Consensus finally...

most calculations confirm the geometrical estimate of cross section;

$$\sigma(XX \to BH) = \pi R_S^2$$

Graviton radiation, etc. that modify the cross sections are included in many new calculations



CMS, JHEP04 (2012)61